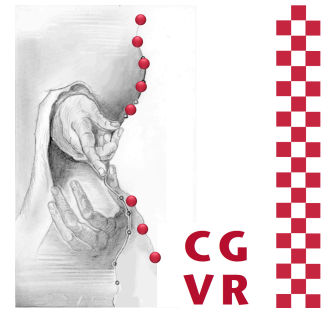
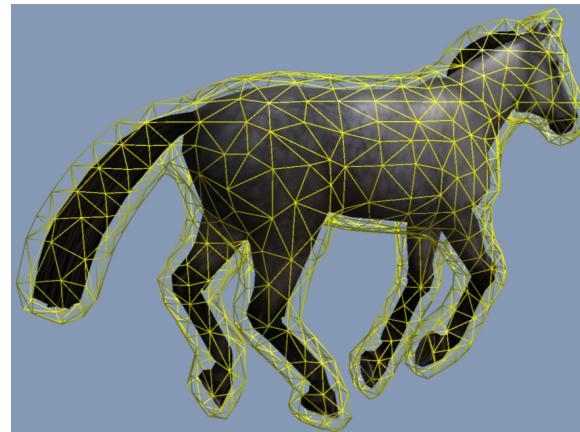


Bremen



Virtual Reality & Physically-Based Simulation Mass-Spring-Systems



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1. Law (law of inertia):

A body, which no forces act upon, continues to move with constant velocity.

- A resting body is just a special case of this law.

2. Law (law of action):

If a force \mathbf{F} acts on a body with mass m , then the body accelerates, and its acceleration is given by

$$\mathbf{F} = m \cdot \mathbf{a}$$

- In other words: force and acceleration are **proportional** to each other; (the proportionality factor happens to be m). In particular, both force and acceleration have the same direction.

3. Law (law of reaction):

If a force \mathbf{F} , that acts on a body, is extended to another body,
Then the opposite force $-\mathbf{F}$ acts on that other body.

- In school, you learn: "action= reaction"

4. Law (law of superposition):

If a number of forces $\mathbf{F}_1, \dots, \mathbf{F}_n$ act on a point or body, then they
can be accumulated by vector addition yielding one resulting
force:

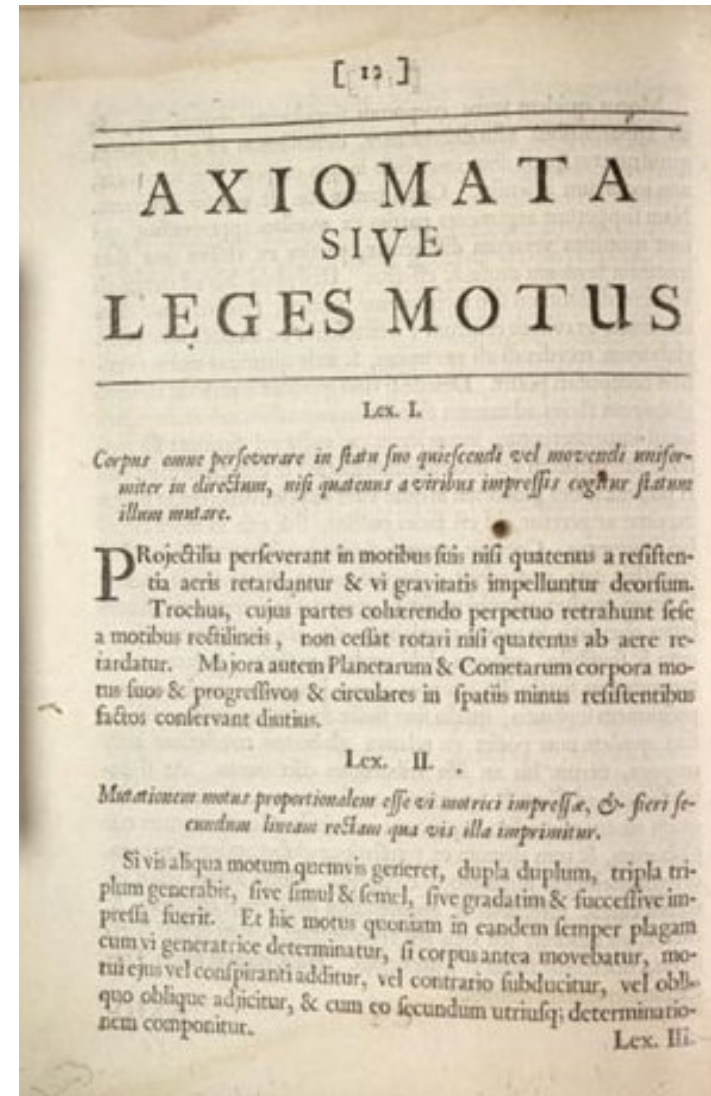
$$\mathbf{F} = \mathbf{F}_1 + \dots + \mathbf{F}_n .$$

- Newton published these laws in his original book

Principia Mathematica

(1687):

- Lex I. Corpus omne perseverare in statu suo quiescendi vel movendi uniformiter in directum, nisi quatenus illud a viribus impressis cogitur statum mutare.*
- Lex II. Mutationem motus proportionalem esse vi motrici impressae, et fieri secundum lineam rectam qua vis illa imprimitur.*



- Definition:

A **spring-mass-system** is a system, consisting of:

1. A set of point masses m_i with positions \mathbf{x}_i and velocities $\mathbf{v}_i, i = 1 \dots N$;
2. A set of springs $s_{ij} = (i, j, k_s, k_d)$, where s_{ij} connects masses i und j , with rest length l_0 , spring constant k_s (= stiffness) and the damping coefficient k_d

- Advantages:

- Very easy to program
- Ideally suited to study different kinds of solving methods
- *Ubiquitous in games* (cloths, capes, sometimes also for deformable objects)

- Disadvantages:

- Some parameters (in particular the spring constants) are **not obvious, i.e.**, difficult to derive
- No volumetric effects (e.g., preservation of volume)

A Single Spring (plus Damper)

- Given: masses m_i and m_j with positions \mathbf{x}_i , \mathbf{x}_j

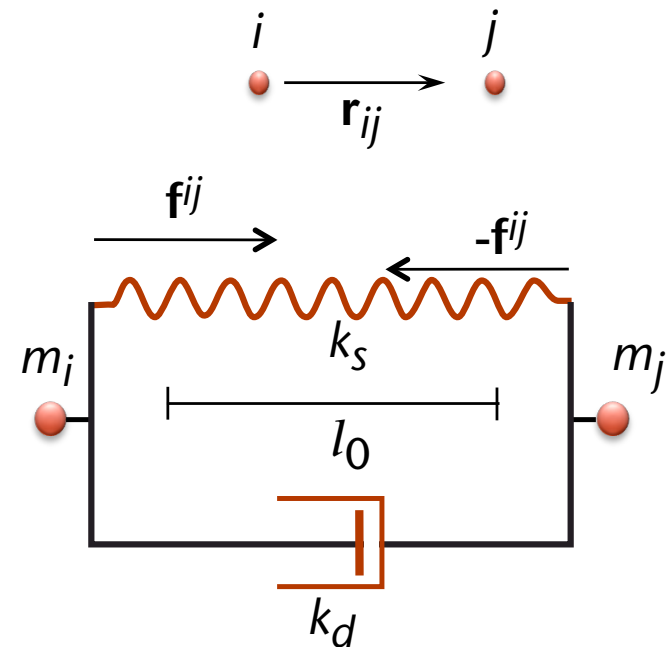
- Let $\mathbf{r}_{ij} = \frac{\mathbf{x}_j - \mathbf{x}_i}{\|\mathbf{x}_j - \mathbf{x}_i\|}$

- The force between particles i and j :

- Force extended by spring:

$$\mathbf{f}_s^{ij} = k_s \mathbf{r}_{ij} (\|\mathbf{x}_j - \mathbf{x}_i\| - l_0)$$

acts on mass m_i in direction of m_j



- Force extended by damper : $\mathbf{f}_d^{ij} = k_d ((\mathbf{v}_j - \mathbf{v}_i) \cdot \mathbf{r}_{ij}) \mathbf{r}_{ij}$

- Sum of forces : $\mathbf{f}^{ij} = \mathbf{f}_s^{ij} + \mathbf{f}_d^{ij}$

- Force on m_j : $\mathbf{f}^{ji} = -\mathbf{f}^{ij}$

- Notice: (4) → the momentum is preserved
 - Momentum = force × mass = $\mathbf{F} \cdot m$
- Note on terminology:
 - German "**Kraftstoß**" = English "**Impulse**" = force × time
 - German "**Impuls**" = English "**momentum**" = force × mass
- Alternative Federkraft: $\mathbf{f}_s^{ij} = k_s \mathbf{r}_{ij} \frac{\|\mathbf{x}_j - \mathbf{x}_i\| - l_0}{l_0}$
- A spring-damper element in reality:



Simulation of a Single Spring

- From Newton's law, we have: $\ddot{\mathbf{x}} = \frac{1}{m} \mathbf{f}$
- Convert differential equation (DE) of order 2 into DE of order 1:

$$\dot{\mathbf{v}}(t) = \frac{1}{m} \mathbf{f}(t)$$

$$\dot{\mathbf{x}}(t) = \mathbf{v}(t)$$

- Initial values (boundary values): $\mathbf{v}(t_0) = \mathbf{v}_0$, $\mathbf{x}(t_0) = \mathbf{x}_0$
- "Simulation" = "Integration of DE's over time"

- By Taylor expansion we get:

$$\mathbf{x}(t + \Delta t) = \mathbf{x}(t) + \Delta t \dot{\mathbf{x}}(t) + O(\Delta t^2)$$

- Analogously: $\mathbf{v}(t + \Delta t) = \mathbf{v}(t) + \Delta t \dot{\mathbf{v}}(t)$

→ This integration scheme is called **explicit Euler integration**


```
forall particles  $i$  :  
    initialize  $\mathbf{x}_i, \mathbf{v}_i, m_i$   
  
Loop forever:  
    forall particles  $i$  :  
         $\mathbf{f}_i \leftarrow \mathbf{f}^g + \mathbf{f}_i^{coll} + \sum_{j, (i,j) \in S} \mathbf{f}(\mathbf{x}_i, \mathbf{v}_i, \mathbf{x}_j, \mathbf{v}_j)$   
  
        forall particles  $i$  :  
             $\mathbf{v}_i += \Delta t \cdot \frac{\mathbf{f}_i}{m_i}$   
  
             $\mathbf{x}_i += \Delta t \cdot \mathbf{v}_i$   
  
    render system every  $n$ -th time
```

\mathbf{f}^g = gravitational force

\mathbf{f}^{coll} = penalty force exerted by collision (e.g., with obstacles)

- Advantages:

- Can be implemented very easily
- Fast execution per time step

- Disadvantages:

- Stable only for very small time steps
 - Typically $\Delta t \approx 10^{-4} \dots 10^{-3}$ sec!
- With large time steps, additional energy is generated "out of thin air", until the system explodes 😊
- Example: overshooting when simulating a single spring
- Errors accumulate quickly

Example for the Instability of Euler Integration

- Consider the differential equation

$$\dot{x}(t) = -kx(t)$$

- The exact solution:

$$x(t) = x_0 e^{-kt}$$

- Euler integration does this:

$$x^{t+1} = x^t + \Delta t(-kx^t)$$

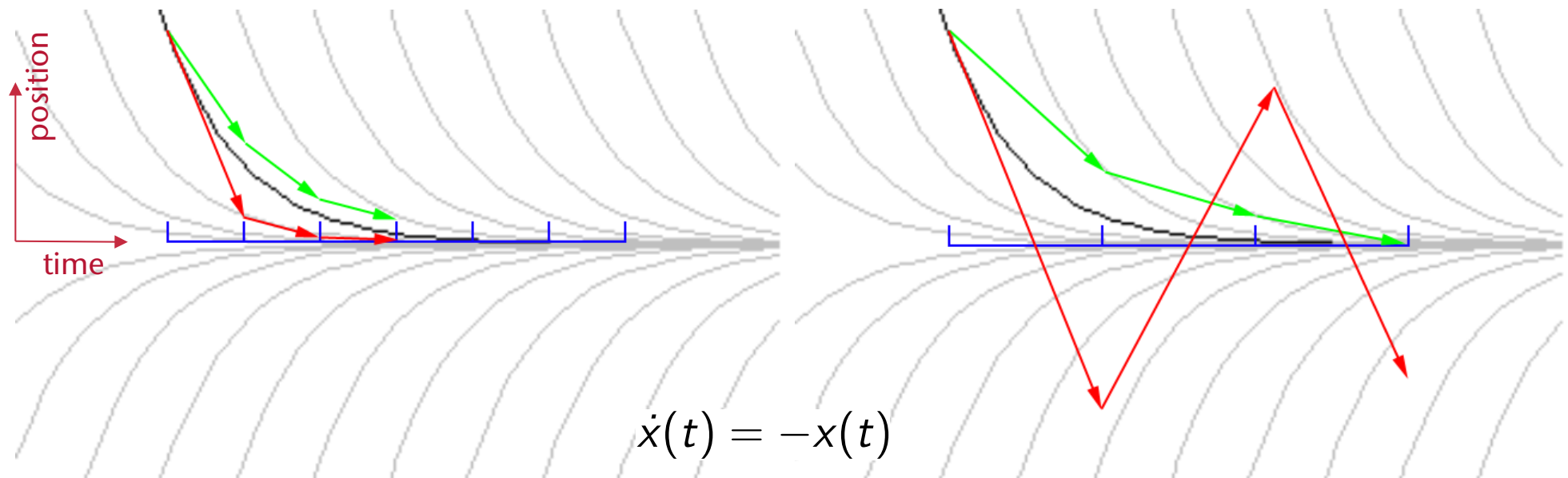
- Case $\Delta t > \frac{1}{k}$:

$$x^{t+1} = x^t \underbrace{(1 - k\Delta t)}_{<0}$$

$\Rightarrow x^t$ oscillates about 0, but approaches 0 (hopefully)

- Case $\Delta t > \frac{2}{k}$: $\Rightarrow x^t \rightarrow \infty$!

■ Visualization:



- Terminology: if k is large \rightarrow the DE is called "*stiff*"
 - The stiffer the DE, the smaller Δt has to be

Visualization of Error Accumulation

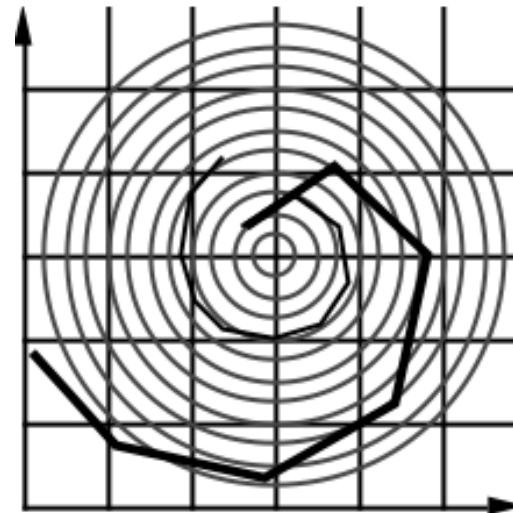
- Consider this DE:

$$\dot{\mathbf{x}}(t) = \begin{pmatrix} -x_2 \\ x_1 \end{pmatrix}$$

- Exact solution:

$$\mathbf{x}(t) = \begin{pmatrix} r \cos(t + \phi) \\ r \sin(t + \phi) \end{pmatrix}$$

- The solution by Euler integration moves in spirals outward, no matter how small Δt !
- Conclusion: Euler integration accumulates errors, no matter how small Δt !

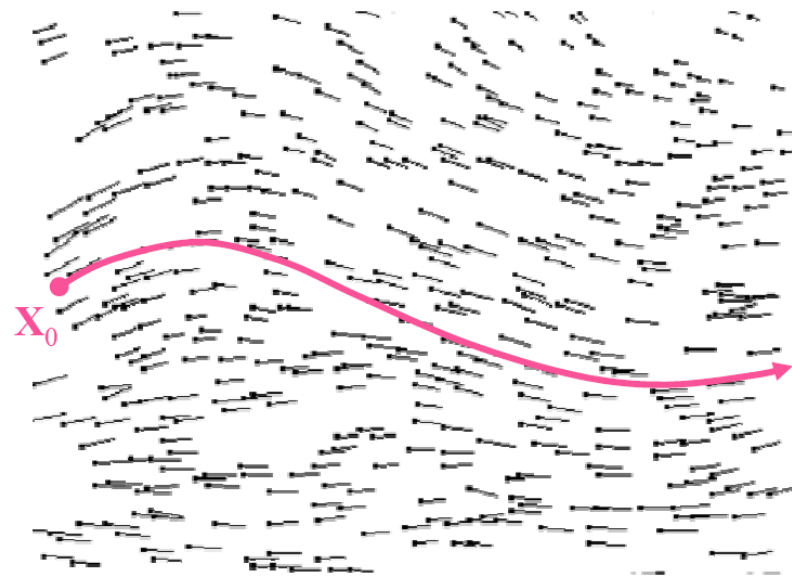


Visualization of Differential Equations

- The general form of a DE:

$$\dot{\mathbf{x}}(t) = \mathbf{f}(\mathbf{x}(t), t)$$

- Visualization of \mathbf{f} as a vector field:



- Notice: this vector field can vary over time!
- Solution of a boundary value problem = path through this field

- Runge-Kutta of order 2:

- Idea: approximate $\mathbf{f}(\mathbf{x}(t), \mathbf{v}(t))$ by a quadratic function that is defined at positions $\mathbf{x}(t)$, $\mathbf{x}(t + \frac{1}{2}\Delta t)$ and $\mathbf{v}(t)$
- The integrator (w/o proof):

$$\mathbf{a}_1 = \mathbf{v}^t \qquad \mathbf{a}_2 = \frac{1}{m} \mathbf{f}(\mathbf{x}^t, \mathbf{v}^t)$$

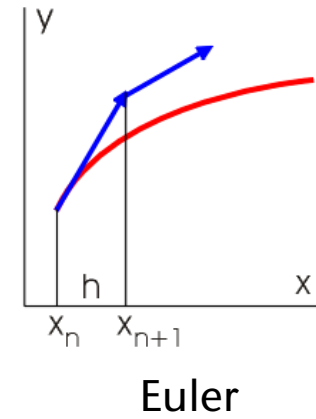
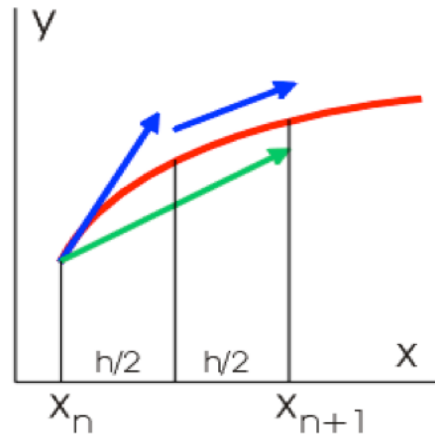
$$\mathbf{b}_1 = \mathbf{v}^t + \frac{1}{2} \Delta t \mathbf{a}_2 \qquad \mathbf{b}_2 = \frac{1}{m} \mathbf{f}\left(\mathbf{x}^t + \frac{1}{2} \Delta t \mathbf{a}_1, \mathbf{v}^t + \frac{1}{2} \Delta t \mathbf{a}_2\right)$$

$$\mathbf{x}^{t+1} = \mathbf{x}^t + \Delta t \mathbf{b}_1 \qquad \mathbf{v}^{t+1} = \mathbf{v}^t + \Delta t \mathbf{b}_2$$

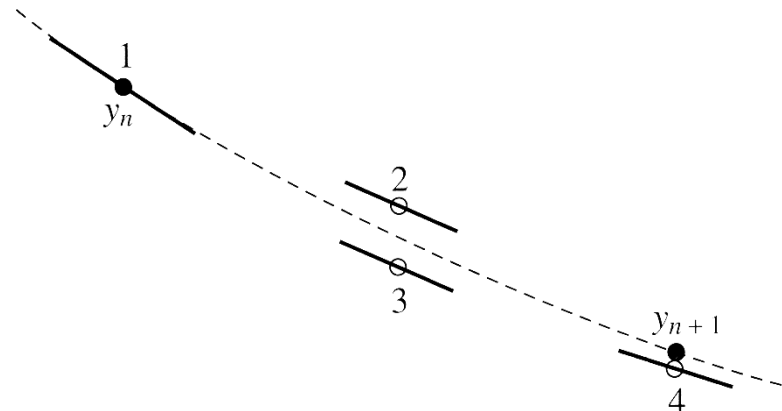
- Runge-Kutta of order 4:

- **The** standard integrator among the explicit integration schemata
- Needs 4 function evaluations (i.e., force computations) per time step
- Order of convergence is: $e(\Delta t) = O(\Delta t^4)$

- Runge-Kutta of order 2:



- Runge-Kutta of order 4:



- A general, alternative method to increase the order of convergence: utilizes values from **history**
- Verlet: utilize $\mathbf{x}(t - \Delta t)$
- Derivation:
 - Develop The taylor series in both time directions:

$$\mathbf{x}(t + \Delta t) = \mathbf{x}(t) + \Delta t \dot{\mathbf{x}}(t) + \frac{1}{2} \Delta t^2 \ddot{\mathbf{x}}(t) + \frac{1}{6} \Delta t^3 \dddot{\mathbf{x}}(t) + O(\Delta t^4)$$

$$\mathbf{x}(t - \Delta t) = \mathbf{x}(t) - \Delta t \dot{\mathbf{x}}(t) + \frac{1}{2} \Delta t^2 \ddot{\mathbf{x}}(t) - \frac{1}{6} \Delta t^3 \dddot{\mathbf{x}}(t) + O(\Delta t^4)$$

- Add both:

$$\mathbf{x}(t + \Delta t) + \mathbf{x}(t - \Delta t) = 2\mathbf{x}(t) + \Delta t^2 \ddot{\mathbf{x}}(t) + O(\Delta t^4)$$

$$\mathbf{x}(t + \Delta t) = 2\mathbf{x}(t) - \mathbf{x}(t - \Delta t) + \Delta t^2 \ddot{\mathbf{x}}(t) + O(\Delta t^4)$$

- Initialization:

$$\mathbf{x}(\Delta t) = \mathbf{x}(0) + \Delta t \mathbf{v}(0) + \frac{1}{2} \Delta t^2 \left(\frac{1}{m} \mathbf{f}(\mathbf{x}(0), \mathbf{v}(0)) \right)$$

- Remark: the velocity does not occur anymore (explicitly)

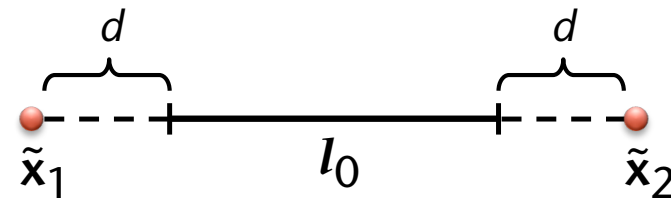
- Big advantage of Verlet over Euler & Runge-Kutta:
it is very easy to handle constraints
- Definition: **Constraint** = some condition on position of one or more mass points
- Examples:
 1. A point must not penetrate an obstacle
 2. The distance between two points must be constant,
or distance must be \leq some specific distance

■ Examples:

- Consider the constraint:

$$\|\mathbf{x}_1 - \mathbf{x}_2\| \stackrel{!}{=} l_0$$

1. Perform one Verlet integration step $\rightarrow \tilde{\mathbf{x}}^{t+1}$
2. Enforce the constraint:



$$\mathbf{x}_1^{t+1} = \tilde{\mathbf{x}}_1^{t+1} + \frac{1}{2} \mathbf{r}_{12} \cdot (\|\tilde{\mathbf{x}}_2^{t+1} - \tilde{\mathbf{x}}_1^{t+1}\| - l_0)$$

$$\mathbf{x}_2^{t+1} = \tilde{\mathbf{x}}_2^{t+1} - \underbrace{\frac{1}{2} \mathbf{r}_{12} \cdot (\|\tilde{\mathbf{x}}_2^{t+1} - \tilde{\mathbf{x}}_1^{t+1}\| - l_0)}_d$$

- Problem: if several constraints are to constrain the same mass point, we need to employ constraint algorithms